



MIGRATE-2017:154246

EVALUATION OF WALL HEAT FLUX BOUNDARY CONDITION IN DSMC IMPLEMENTED IN OPENFOAM

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KEYWORDS

Wall heat flux boundary condition, Nanochannel, Nanocavity, DSMC.

ABSTRACT

Cooling and heating play a major role in micro/nano-devices and rarefied gases [1]. In this paper, we report the validation of a wall heat flux boundary condition, implemented in the direct simulation Monte Carlo (DSMC) solver of the OpenFOAM package. We used an enhanced version of OpenFOAM 2.4.0 called OpenFOAM 2.4-MNF. The DSMC solver in this release has enhanced features such as transient adaptive subcells, chemical reaction modules, and defined pressure inlet/outlet boundary conditions. The modified iterative technique (MIT), which was recently introduced by Akhlaghi and Roohi [2] was used in this implementation. This approach implements a desired wall heat flux distribution over the walls of the geometry for rarefied gas simulations.

In this technique, the local wall temperature is modified iteratively in such a way that the desired wall heat flux is achieved. The wall temperature is updated using:

$$\Delta T_w = \alpha (q_w^* - q_{des}^*) T_w, \quad (1)$$

where α is a parameter controlling convergence, T_w is the current wall temperature, ΔT_w is the change in wall temperature required to iterate towards the desired heat flux, and q_{des}^* and q_w^* denote the desired and obtained local wall heat flux, respectively. The heat fluxes are normalized as follows:

$$q_w^* = \frac{q_w}{q_f} \quad (2)$$

where q_f is a macroscopic characteristic heat flux defined as

$$q_f = C_p \rho_0 U_0 T_0$$

where C_p , ρ_0 , U_0 , and T_0 are reference values of specific heat capacity, density, velocity and temperature, respectively. (3)

Here, we validate the implementation of equation (1) in dsmcFoam. The working gas is argon in a 2D pressure-driven nanochannel. Fig. 1 shows the normalized center-line pressure distribution under three

different wall boundary conditions; cooling, adiabatic, and heating wall boundary conditions. Our dsmcFoam solutions are compared with the DSMC solution of Balaj et al. [3] under the same wall conditions and good agreement is found. As is shown, heating or cooling affects the pressure distribution. Cooling results in a concave pressure distribution, and heating increases the convex nature of the pressure profile.

Fig. 2 demonstrates the effect of the heating condition on the wall temperature distribution of the nanocavity. The Knudsen number, Kn , is 0.1, the top wall is set at a constant temperature of 500 K, and the three other walls at a constant heat flux of $q^*=-0.1$. The right frame in the figure shows temporal variations of the wall temperature and matches the DSMC solution reported in [2]. Consequently, results shown in Figs. 1-2 indicate that our implementation of the MIT technique is accurate. The next stage will consider wall heat flux effects on the flow and thermal field behavior in complex geometries.

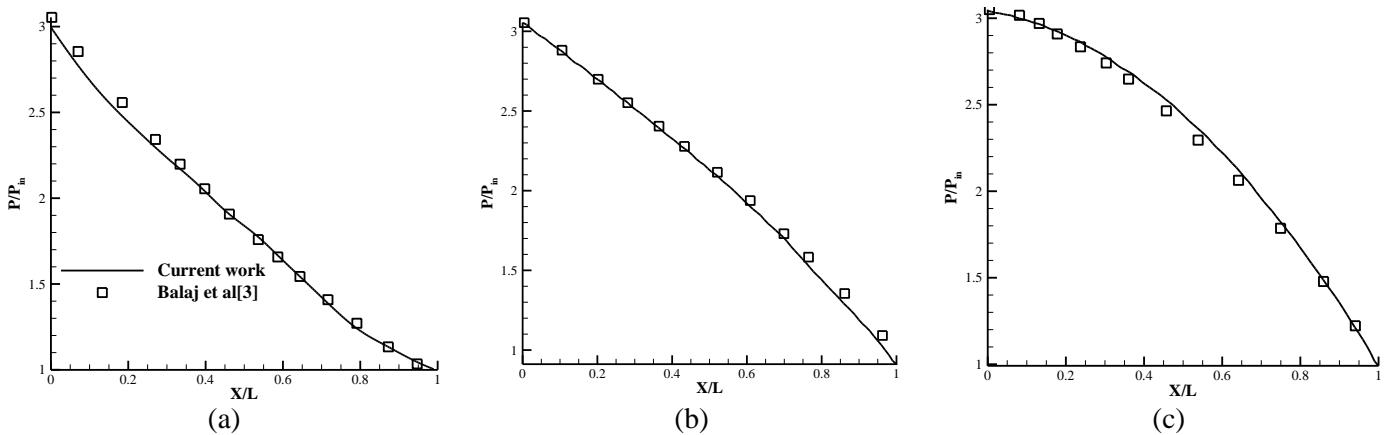


Figure 1: Flow in a microchannel a) Cooling condition ($Kn=0.04$, $q^*=-0.3$), b) Adiabatic ($Kn=0.03$, $q^*=0$), and c) Heating ($Kn=0.02$, $q^*=+0.3$)

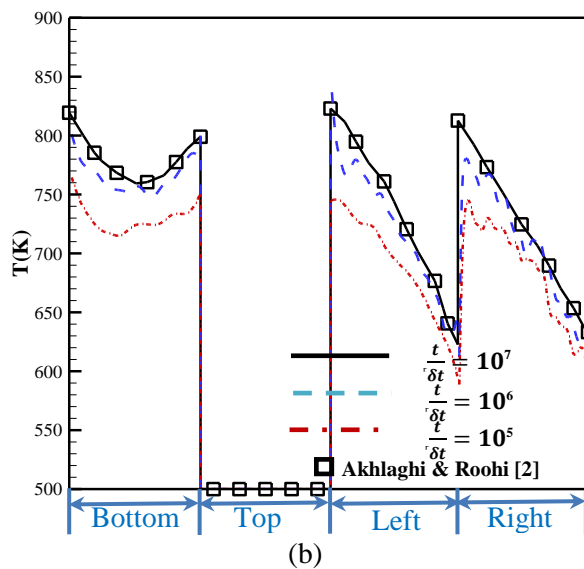
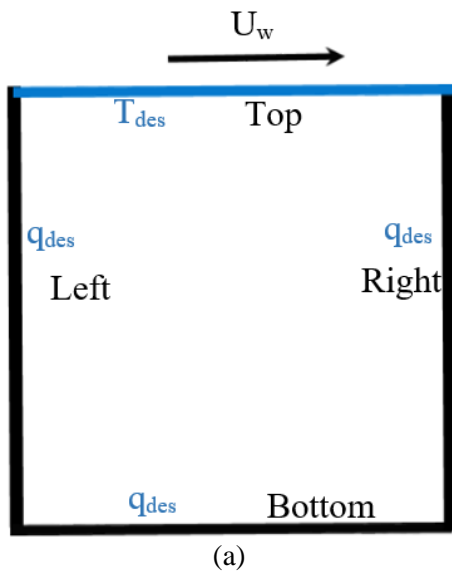


Figure 2: a) Schematic geometry of cavity, and b) Temperature profile over four sides of the cavity, $Kn=0.1$, $q^*=-0.1$.

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within H2020



Proceedings of the 2nd MIGRATE Workshop
June 29-30, 2017 – Sofia, Bulgaria

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